Trigonometric Function Solutions of Fractional Drinfeld’s Sokolov -Wilson System

Zeynep Fidan Koçak and Gülnur Yel

1Department of Mathematics, University of Mugla Sıtkı Koçman, Mugla, Turkey
2Department of Mathematics Education, Final International University, Kyrenia, Cyprus

Abstract: In this paper, we construct exact trigonometric solutions of the space-time fractional classical Drinfeld’s Sokolov-Wilson system by Modified Trial Equation Method (MTEM). These solutions may explain some physical phenomena and lead to researchers in physics and engineering.

1. Introduction

Studies in recent years on fractional analysis have been increasingly proceed. Fractional analysis describe many scientific phenomena such as physics, chemistry, biology, engineering and so on. Several efficient analytical, semi-analytical and numerical methods are applied for solving fractional equations and equation systems. Some of them are Modified Trial Equation method, Sumudu Transform method, Riccati-Bernoulli sub-ODE method, Improved Bernoulli sub-equation method, modified Kudryashov method, Functional Variable method, Homotopy Analysis method and so on have been used to find new solutions of fractional equations [1-11].

We apply Modified Trial Equation method (MTEM) for obtain new travelling wave solution to fractional Drinfeld’s Sokolov-Wilson system [1,12].

2. The Modified Trial Equation Method

In this subsection we describe structure of the MTEM [2-5].

We consider partial differential equation in two variables and a dependent variable \( u \):

\[
P(u,u_x,u_{xx},u_{xxx},...) = 0, \quad (2.1)
\]

and take the wave transformation,

\[
u(x,t) = u(\xi), \xi = kx - ct, \quad (2.2)
\]

where \( k \) and \( c \) are constants can be determined later. By substituting Eq. (2.2) into Eq. (2.1), a nonlinear ordinary differential equation (NODE) is converted as following:

\[
N(U, U', U'', U''', ...) = 0. \quad (2.3)
\]

We assume that the solution can be expressed in the form

\[
U' = \frac{F(u)}{G(u)} = \frac{\sum_{i=0}^{v} a_i u^i}{\sum_{j=0}^{l} b_j u^j} = \frac{a_0 + a_1 u + a_2 u^2 + \cdots + a_v u^v}{b_0 + b_1 u + b_2 u^2 + \cdots + b_l u^l}, \quad (2.4)
\]

\[
U^* = \frac{F(u)}{G(u)} \left( F'(u) G(u) - F(u) G'(u) \right) \frac{1}{G'(u)}, \quad (2.5)
\]

where \( F(u) \) and \( G(u) \) are polynomials. Substituting above relations into Eq.(2.3) yields an equation of polynomial \( \Omega(u) \) of \( u \):

\[
\Omega(u) = \rho_1 u' + \cdots + \rho_s u + \rho_0 = 0. \quad (2.6)
\]

According to the balance principle, we can get a relationship between \( n \) and \( l \). We can compute some values of \( n \) and \( l \).

Let the coefficients of \( \Omega(u) \) all be zero will yield an algebraic equations system:

\[
\rho_i = 0, i = 0,1,2,\cdots, s. \quad (2.7)
\]

By solving this system, we will thus determine the values of \( a_0, a_1,\ldots, a_v \) and \( b_0, b_1,\ldots, b_l \).

Reduce Eq.(2.4) to the elementary integral form,

\[
\pm (\mu - \mu_0) = \int \frac{G(u)}{F(u)} du. \quad (2.8)
\]
Using a complete discrimination system for polynomial of \( F(u) \), we solve Eq.(2.8) with the help of Mathematica 9 and classify the exact solutions to Eq.(2.3). For better explication of results obtained in this way, we can plot two and three dimensional surfaces of the solutions obtained by using suitable parameters.

3. Application

Let consider the the space-time fractional classical Drinfeld's Sokolov-Wilson system as follows [1],

\[
\frac{\partial^\alpha u}{\partial t^\alpha} + p\frac{\partial^\beta v}{\partial x^\beta} = 0,
\]

\[
\frac{\partial^\alpha v}{\partial t^\alpha} + q\frac{\partial^\beta v}{\partial x^\beta} + ru\frac{\partial^\beta v}{\partial x^\beta} + sv\frac{\partial^\beta u}{\partial x^\beta} = 0, \tag{3.1}
\]

where \( u \) and \( v \) are the functions of \((x,t)\), \(0 < \alpha, \beta \leq 1, x > 0\). We apply following transformations then the Eq. (3.1) can be reduced to ordinary differential equation.

\[
u(x,t) = U(\eta), v(x,t) = V(\eta), \eta = \frac{wx^\alpha}{\Gamma(1+\beta)} - \frac{\lambda t^\alpha}{\Gamma(1+\alpha)}, \tag{3.2}
\]

where \( w \) and \( \lambda \) are arbitrary constants and \( a \neq 1, b \neq 1 \) are fractal constants. We get differentiate the function \( U(\eta), V(\eta) \) respect to \( \eta \), then Eq.(3.2) is rearranged, it yields us,

\[-\lambda bU' + pawVV' = 0, \tag{3.3}
\]

\[-\lambda bV' + qa^3w^3V''' + rawUV' + sawVU' = 0. \tag{3.4}
\]

We obtain following equalities from Eq.(3.3),

\[U' = \left(\frac{paw}{b\lambda}\right)VV', \tag{3.5}
\]

\[U = \left(\frac{paw}{b\lambda}\right)V^2 + c_1, \tag{3.6}
\]

where \( c_1 \) is constant of integration. Embedding Eqs.(3.5,3.6) into Eq.(3.4), we have the nonlinear ordinary differential equation:

\[qa^3w^3V''' + pa^2w^2\left(\frac{r + 2s}{2b\lambda}\right)V^2V' + (rawc_1 - b\lambda)V' = 0. \tag{3.7}
\]

Integrating Eq.(3.7) once,

\[qa^3w^3V'' + pa^2w^2\left(\frac{r + 2s}{2b\lambda}\right)V^2V' + (rawc_1 - b\lambda)V' = 0. \tag{3.8}
\]

where \( c_1 \) is constant of integration. When we reconsider the Eq.(3.8) for homogenous balance principle between \( V' \) and \( V^3 \), we obtain the following relationship for \( n \) and \( l \),

\[2n - 2l - 1 = 3 \Rightarrow n = l + 2, \tag{3.9}
\]

Case 1: For the values of \( l = 0, n = 2 \), we get;

\[V' = F(\nu) = \frac{\sum a_i^jv^i}{b_0}, \tag{3.10}
\]

\[V^* = \frac{F'(\nu)F(\nu) - F(\nu)G'(\nu)}{G'(\nu)} = \frac{(a_i + 2a_i v)(a_i + a_i v + a_i v^2)}{b_0^3}, \tag{3.11}
\]

where \( a_2 \neq 0 \) and \( b_0 \neq 0 \). Using the Wolfram Mathematica 9, the algebraic equation system is solved,

\[a_1 = 0, c_2 = 0, r = \frac{\lambda bb_0^2 - 2qa^3w^3a_i a_2}{aw b_0 c_i}, \tag{3.12}
\]

\[p = -\frac{12a^2bqw^2\lambda a_2 c_i}{2qa^3w^3a_0 - b_0^2 (b\lambda + 2awc_1)}.
\]

we have coefficients above. By substituting these coefficients in Eq.(3.8), we obtain the solution of \( v(x,t) \) and \( u(x,t) \)

\[v(x,t) = \sqrt{a_2}\tan\left(a_2^\alpha\frac{wx^\alpha}{b_0} - \frac{\lambda t^\alpha}{\Gamma(1+\alpha)}\right)^{-\frac{1}{2}}, \tag{3.13}
\]

\[u(x,t) = \left(a_2^\alpha\frac{wx^\alpha}{b_0} - \frac{\lambda t^\alpha}{\Gamma(1+\alpha)}\right)^{\frac{1}{2}} + c_i \tag{3.14}
\]

* Corresponding author: gulnurvel33@gmail.com
it yields us, ordinary differential equation: Eqs.(3.5,3.6) into Eq.(3.4), we have the nonlinear are fractal con stants. We get differentiate the function

\[ \frac{\partial^\alpha v}{\partial x^\alpha} + \frac{\partial^\beta v}{\partial y^\beta} = f(x,y) \]

where \( a, b \) are arbitrary constants and \( \alpha, \beta \) are fractal con stants, we can say that this method is applied to nonlinear differential equations and systems. These trigonometric function solutions have been introduced to the literature for the first time. We think that these new solutions lead the way other scientific area.

4. Conclusions

In this manuscript, we have efficiently and easily practice the Modified Trial Equation Method that is give analytical solution. We can say that this method is applied to nonlinear differential equations and systems. These trigonometric function solutions have been introduced to the literature for the first time. We think that these new solutions lead the way other scientific area.

References


* Corresponding author: gulnuryel33@gmail.com

Fig. 1. The 2D and 3D surfaces for the the space-time fractional classical Drinfeld's Sokolov-Wilson system with

\[ b_0 = 0.2, \lambda = 0.3, a = b = 0.1, \alpha = \beta = 0.5, \]

\[ c_1 = 1, w = 0.9, s = 0.04, 0 < x < 20, 0 < t < 10, \]

\[ q = 0.04, a_x = 0.1, a_y = 0.9, \text{ and } t = 0.4, \text{for 2D surfaces.} \]